# Finite Difference Time Domain Application to room acoustics 

## Master 2 Acoustical Engineering

 Numerical Techniques for Acoustics - Session 4
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## Wave equation in time domain

For a wave on a vibrating string, magnitude of the vibration $y(x, t)$ is governed by a Partial Derivative Equation (PDE) :

$$
\begin{aligned}
\frac{1}{c^{2}} \frac{\partial^{2} y}{\partial t^{2}}-\frac{\partial^{2} y}{\partial x^{2}} & =0 \\
y(x=0, \cdot) & =0 \\
y(x=L, \cdot) & =0 \\
y(\cdot, t=0) & =y_{0} \\
\frac{\partial y}{\partial t}(\cdot, t=0) & =v_{0}
\end{aligned}
$$

with $c$ the sound celerity, $t$ the time, $x$ the position on the string, $L$ the length of the string and $y_{0}, v_{0}$ the initial position and speed.

Note : 4 derivatives $\rightarrow 4$ initial conditions.

## Classical scheme for space derivative

As usual, we consider Taylor's expansion for all functions $y \in C^{4}([0, L])$ :

$$
y^{\prime \prime}(x)=\frac{y(x+h)-2 y(x)+y(x-h)}{h^{2}}+O\left(h^{4}\right)
$$

Considering a string discretization by a regular grid $X=\left(x_{0}, \ldots, x_{N}\right)$, such as :

- $x_{0}=0$,
- $x_{N}=L$,
- $x_{n}=n \delta_{x}$ with $\delta_{x}=\frac{L}{N+1}$ and $n \in[0, N]$,
a centered scheme can be used :

$$
\begin{aligned}
y^{\prime \prime}(x) & =\frac{y\left(x+\delta_{x}\right)-2 y(x)+y\left(x-\delta_{x}\right)}{\delta_{x}^{2}}+O\left(\delta_{x}^{4}\right) \\
y_{n}^{\prime \prime} & \approx \frac{y_{n+1}-2 y_{n}+y_{n-1}}{\delta_{x}^{2}}
\end{aligned}
$$

## Classical scheme for time derivative

Fixing a final time $t_{f}$ for the wave propagation, a time discretisation $T=\left(t_{0}, \ldots, t_{P}\right)$ can be defined such as:

- $t_{0}=0$,
- $t_{P}=t_{f}$,
- $t_{p}=p \delta_{t}$ with $\delta_{t}=\frac{t_{f}}{P+1}$ and $p \in[0, P]$.

In the same way as for space, second order derivative in time leads to another centered sheme :

$$
\begin{aligned}
y^{\prime \prime}(t) & =\frac{y\left(t+\delta_{t}\right)-2 y(t)+y\left(t-\delta_{t}\right)}{\delta_{t}^{2}}+O\left(\delta_{t}^{4}\right) \\
y_{p}^{\prime \prime} & \approx \frac{y_{p+1}-2 y_{p}+y_{p-1}}{\delta_{t}^{2}}
\end{aligned}
$$

Note : Not so far from the space scheme...

## Classical scheme for wave equation

Combine each partial derivative approximation leads to the leap-frog scheme:

$$
\begin{gathered}
\frac{1}{c^{2}} \frac{\partial^{2} y}{\partial t^{2}}-\frac{\partial^{2} y}{\partial x^{2}}=0 \\
\frac{1}{c^{2}} \frac{y_{n}^{p+1}-2 y_{n}^{p}+y_{n}^{p-1}}{\delta_{t}^{2}}-\frac{y_{n+1}^{p}-2 y_{n}^{p}+y_{n-1}^{p}}{\delta_{x}^{2}}=0
\end{gathered}
$$

- Time indices up, space indices down,
- Implicit scheme in space,
- Explicit scheme in time (two step scheme),
- Stable $\Longleftrightarrow \frac{\delta_{x}}{\delta_{t}} \geq c$ (CFL condition) .

Note: CFL from Richard Courant, Kurt Friedrichs and Hans Lewy : the "speed" of the scheme has to be greater than the speed of the equation.

## Algorithm

- Use sparse matrix for the discrete derivative of the space part:

$$
M_{x}=\frac{1}{\delta_{x}^{2}}\left(\begin{array}{cccccc}
-2 & 1 & 0 & \ldots & 0 & 0 \\
1 & -2 & 1 & \ldots & 0 & 0 \\
0 & 1 & -2 & \ldots & 0 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
0 & 0 & 0 & \ldots & -2 & 1 \\
0 & 0 & 0 & \ldots & 1 & -2
\end{array}\right)
$$

- Add initial condition to the linear system :

$$
M_{x}=\frac{1}{\delta_{x}^{2}}\left(\begin{array}{cccccc}
\delta_{x}^{2} & 0 & 0 & \ldots & 0 & 0 \\
1 & -2 & 1 & \ldots & 0 & 0 \\
0 & 1 & -2 & \ldots & 0 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
0 & 0 & 0 & \ldots & -2 & 1 \\
0 & 0 & 0 & \ldots & 0 & \delta_{x}^{2}
\end{array}\right)
$$

## Algorithm

- Initialize 2 vectors of unknowns in space:

$$
\begin{gathered}
Y^{0}=\left(\begin{array}{c}
y_{0}^{0} \\
y_{1}^{0} \\
\ldots \\
y_{N}^{0}
\end{array}\right)=\left(\begin{array}{c}
y\left(x_{0}, 0\right) \\
y\left(x_{1}, 0\right) \\
\ldots \\
y\left(x_{N}, 0\right)
\end{array}\right) \\
Y^{1}=\left(\begin{array}{c}
y_{0}^{1} \\
y_{1}^{1} \\
\ldots \\
y_{N}^{1}
\end{array}\right)=Y^{0}+\delta_{t}\left(\begin{array}{c}
\partial_{t} y\left(x_{0}, 0\right) \\
\partial_{t} y\left(x_{1}, 0\right) \\
\ldots \\
\partial_{t} y\left(x_{N}, 0\right)
\end{array}\right)
\end{gathered}
$$

- Compute $Y^{2}$, using the leap-frog scheme:

$$
Y^{2}=2 Y^{1}-Y^{0}+\left(c \delta_{t}\right)^{2} M x Y^{1}
$$

- Make a recursion until the final time $t_{f}$.


## Starting code

```
% Clean up
clear all
close all
clc
% Physical parameters
L = 1; % String size
x0 = 0.3; % Initial position
tf = 1; % Final time
c = 1; % Sound celerity
```

\% Initial condition for magnitude and speed
$u 0=@(x) \exp (-(x-x 0) . \wedge 2 / 1 e-2) ; \%$ gaussian
$\mathrm{v} 0=$ @ $(x)$ zeros (size (x)); $\%$ null
\% Numerical discretization
$\mathrm{dx}=0.01 ; \quad$ \% Space step
$\mathrm{dt}=0.5 * \mathrm{dx} / \mathrm{c} ; \quad$ \% Time step (following CFL)

## Numerical result



## 1-D air propagation



## 1-D room acoustic

For a wave propagating in a 1-D room, relative magnitude of the wave $u(x, t)$ is governed by a PDE:

$$
\begin{aligned}
\frac{1}{c^{2}} \frac{\partial^{2} u}{\partial t^{2}}-\frac{\partial^{2} u}{\partial x^{2}} & =0 \\
\frac{\partial u}{\partial x}(x=0, \cdot) & =0 \\
\frac{\partial u}{\partial x}(x=L, \cdot) & =0 \\
u(\cdot, t=0) & =u_{0} \\
\frac{\partial u}{\partial t}(\cdot, t=0) & =v_{0}
\end{aligned}
$$

with $c$ the sound celerity, $t$ the time, $x$ the position in the room, $L$ the length of the room and $u_{0}, v_{0}$ the initial position and speed.

Note : Physics is completely different, but equation is almost the same. Only initial conditions in space have changed.

## Algorithm modifications

As only initial conditions in space has changed, only sparse matrix has to be modified :

$$
\begin{aligned}
& \frac{\partial u}{\partial x}(x=0, \cdot)=0 \Longleftrightarrow \frac{u_{0}-u_{1}}{\delta_{x}}=0 \\
& \frac{\partial u}{\partial x}(x=L, \cdot)=0 \Longleftrightarrow \frac{u_{N-1}-u_{N}}{\delta_{x}}=0
\end{aligned}
$$

which implies:

$$
M_{x}=\frac{1}{\delta_{x}^{2}}\left(\begin{array}{cccccc}
-1 & 1 & 0 & \ldots & 0 & 0 \\
1 & -2 & 1 & \ldots & 0 & 0 \\
0 & 1 & -2 & \ldots & 0 & 0 \\
\ldots & \ldots & \ldots & \ldots & \ldots & \ldots \\
0 & 0 & 0 & \ldots & -2 & 1 \\
0 & 0 & 0 & \ldots & 1 & -1
\end{array}\right)
$$

Note : $u(x=0, \cdot)=0$ is a Dirichlet condition, $\frac{\partial u}{\partial x}(x=0, \cdot)=0$ is a Neumann condition.

## Numerical result



## 2-D air propagation



2-D bounded domain


## 2-D room acoustic

For a wave propagating in a 2-D room $\Omega^{i}$, relative magnitude of the wave $u(\mathbf{x}, t)=u(x, y, t)$ is governed by a PDE :

$$
\begin{aligned}
\frac{1}{c^{2}} \frac{\partial^{2} u}{\partial t^{2}}-\Delta_{\mathrm{x}} u & =0 \\
\frac{\partial u}{\partial n}(\mathbf{x} \in \Gamma, \cdot) & =0 \\
u(\cdot, t=0) & =u_{0} \\
\frac{\partial u}{\partial t}(\cdot, t=0) & =v_{0}
\end{aligned}
$$

with $c$ the sound celerity, $t$ the time, $\mathbf{x}=(x, y)$ the position in the room and $u_{0}, v_{0}$ the initial position and speed.
Reminder :

$$
\Delta_{\mathrm{x}} u=\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}} .
$$

Note : Physics is completely different, but equation is almost the same. Only initial conditions in space have changed!

## Still the same approach...

Considering a tensor product with two regular grids :

- $x_{m}=m \delta_{\mathbf{x}}$ for all $m \in[0, M]$,
- $y_{n}=n \delta_{\mathrm{x}}$ for all $n \in[0, N]$,

Taloyr expansion can be applied simultaneously :

$$
\begin{aligned}
\Delta_{\mathbf{x}} u(x, y)= & \frac{\partial^{2} u}{\partial x^{2}}(x, y)+\frac{\partial^{2} u}{\partial y^{2}}(x, y) \\
\Delta_{\mathbf{x}} u(x, y)= & \frac{u(x+h, y)-2 u(x, y)+u(x-h, y)}{h^{2}}+ \\
& \frac{u(x, y+h)-2 u(x, y)+u(x, y-h)}{h^{2}}+O\left(h^{4}\right) \\
\Delta_{\mathbf{x}} u_{m, n} \approx & \frac{u_{m+1, n}+u_{m, n+1}-4 u_{m, n}+u_{m-1, n}+u_{m, n-1}}{\delta_{\mathbf{x}}^{2}}
\end{aligned}
$$

Note: It's a five points sheme (e.g. blackboard).

## Algorithm

- Use sparse matrix for the discrete derivative of the Laplacian,
- Add initial condition (Dirichlet or Neumann)


Example with $-\delta_{\mathbf{x}}^{2} \Delta_{\mathbf{x}}$ matrix with Dirichlet condition.

## Algorithm

- Initialize 2 vectors of unknowns in space :

$$
U^{0}=\left(\begin{array}{c}
u_{0,0}^{0} \\
u_{1,0}^{0} \\
\ldots \\
u_{M, 0}^{0} \\
u_{0,1}^{0} \\
u_{1,1}^{0} \\
\ldots \\
u_{M, N}^{0}
\end{array}\right)=\left(\begin{array}{c}
u\left(x_{0}, y_{0}, 0\right) \\
u\left(x_{1}, y_{0}, 0\right) \\
\ldots \\
y\left(x_{M}, y_{0}, 0\right) \\
u\left(x_{0}, y_{1}, 0\right) \\
u\left(x_{1}, y_{1}, 0\right) \\
\ldots \\
y\left(x_{M}, y_{N}, 0\right)
\end{array}\right)
$$

and $U^{1}=U^{0}+\delta_{t} v_{0}$.

- Compute $Y^{2}$, using the leap-frog scheme :

$$
Y^{2}=2 Y^{1}-Y^{0}+\left(c \delta_{t}\right)^{2} M_{x y} Y^{1}
$$

- Make a recursion until the final time $t_{f}$.


## Starting code

```
% Clean up
clear all
close all
clc
```

\% Physical parameters
$\mathrm{L}=\left[\begin{array}{ll}3 & 2\end{array}\right] ; \quad$ \% Room size
XX $=\left[\begin{array}{ll}2.2 & 1.2\end{array}\right] ; \quad \%$ Initial position
$\mathrm{tf}=5 ; \quad$ \% Final time
c $=1 ; \quad$ \% Sound celerity
\% Initial condition for magnitude and speed
wi $=@(x, y) \exp (-(x-X 0(1)) \cdot \wedge 2 / 1 e-2) \cdot * \ldots$
$\exp (-(y-X 0(2)) . \wedge 2 / 1 e-2) ; \quad \%$ gaussian
vi $=$ @(x,y) zeros(size(x));
\% null
\% Numerical discretization

```
dx = 0.01; % Space step (both for }x\mathrm{ and y)
dt = 0.5*dx/c;
    % Time step (CFL)
```

Numerical result


