

# Feedback generation of quantum entangled states

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# Presentation outline

- 1 Feedback stabilization of single spin systems
- 2 Entanglement
- 3 Symmetric and anti-symmetric entangled states
- 4 General entangled states

# Experience scheme

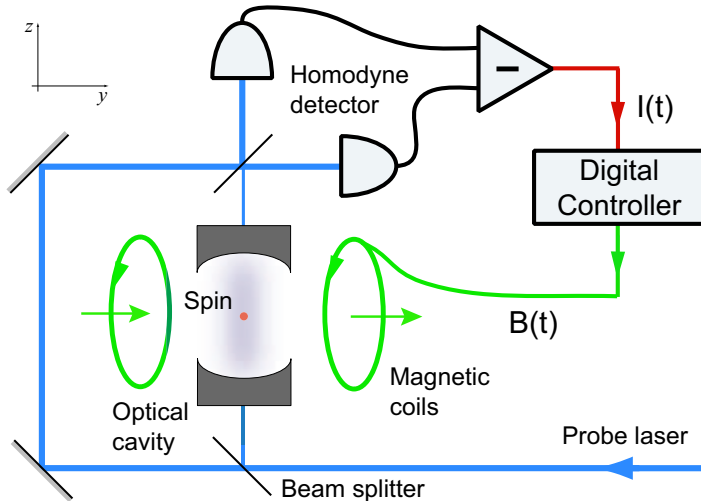


Figure: R. van Handel, J.K. Stockton, H. Mabuchi, IEEE Trans-AC, 2005

# Quantum filter equation

**State space:**  $\{\rho \in \mathbb{C}^{2 \times 2} \mid \rho \geq 0, \text{ Hermitian}, \text{Tr}(\rho) = 1\}$ .

$$d\rho_t = -iu_t[\sigma_y, \rho_t]dt + \frac{1}{2}(2\sigma_z\rho_t\sigma_z - \sigma_z^2\rho_t - \rho_t\sigma_z^2)dt \\ + (\sigma_z\rho_t + \rho_t\sigma_z - 2\text{Tr}(\sigma_z\rho_t)\rho_t)dW_t, \\ dW_t = dy_t - 2\text{Tr}(\sigma_z\rho_t)dt$$

where

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

# Stabilization problem

$$d\rho_t = \frac{1}{2}(2\sigma_z\rho_t\sigma_z - \sigma_z^2\rho_t - \rho_t\sigma_z^2)dt + (\sigma_z\rho_t + \rho_t\sigma_z - 2\text{Tr}(\sigma_z\rho_t)\rho_t)dW_t,$$

## Equilibrium states

$$\rho_{\uparrow} := \Pi_{|\uparrow\rangle} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}, \quad \rho_{\downarrow} := \Pi_{|\downarrow\rangle} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}.$$

## Problem

To stabilize deterministically one of these states.

# Stochastic Lyapunov theory

## Stochastic system:

$$dx_t = f(x_t)dt + \sigma(x_t)dw_t, \quad x(t=0) = x_0,$$

## Lyapunov function:

$$\frac{d}{dt} \mathbb{E} V(x_t) = \mathcal{L}V(x_t) = \frac{\partial V}{\partial x}(x_t)f(x_t) + \frac{1}{2} \frac{\partial^2 V}{\partial x^2}(x_t)\sigma(x_t)^2 \leq 0.$$

- 1 **Doob's inequality:**  $\mathbb{P}(\sup_{0 \leq t < \infty} V(x_t) \geq \alpha) \leq \frac{V(x_0)}{\alpha}$ .
- 2 Convergence in probability towards the invariant set included in  $\mathcal{L}V = 0$  (**Kushner's theorem**).

## Spin-1/2 case

$$d\rho_t = -iu_t[\sigma_y, \rho_t]dt + \frac{1}{2}(2\sigma_z\rho_t\sigma_z - \sigma_z^2\rho_t - \rho_t\sigma_z^2)dt \\ + (\sigma_z\rho_t + \rho_t\sigma_z - 2\text{Tr}(\sigma_z\rho_t)\rho_t)dW_t.$$

**Lyapunov function:**

$$V(\rho_t) = 1 - \text{Tr}(\rho_t\rho_\uparrow)^2$$

$$\mathcal{L}V(\rho_t) = 2u_t\text{Tr}(i[\sigma_y, \rho_t]\rho_\uparrow) - 4\text{Tr}(\rho_t\rho_\uparrow)^2(1 - \text{Tr}(\sigma_z\rho_t))^2.$$

**A possible approach:**

$$u_t = -\text{Tr}(i[\sigma_y, \rho_t]\rho_\uparrow)$$

## Spin-1/2 case (continued)

### Invariant set

$\rho_t$  converges in probability towards the set  $\{\rho_\uparrow, \rho_\downarrow\}$ . The Lyapunov function  $V$  takes its maximal value on  $\rho_\downarrow$ :  $V(\rho_\downarrow) = 1 - \text{Tr}(\rho_\uparrow \rho_\downarrow)^2 = 1$ .

Contrarily to the deterministic case, we do not have semi-global stabilization.

### Change of strategy

- 1  $u_t = -\text{Tr}(i[\sigma_y, \rho_t]\rho_\uparrow)$  if  $\text{Tr}(\rho_t \rho_\uparrow) \geq \gamma$ ;
- 2  $u_t = 1$  if  $\text{Tr}(\rho_t \rho_\uparrow) \leq \gamma/2$ ;
- 3 if  $\rho_t \in \mathcal{B} = \{\rho : \gamma/2 < \text{Tr}(\rho \rho_\uparrow) < \gamma\}$ , then  $u_t = -\text{Tr}(i[\sigma_y, \rho_t]\rho_\uparrow)$  if the last entry of  $\rho_t$  into  $\mathcal{B}$  has been via the boundary  $\text{Tr}(\rho \rho_\uparrow) = \gamma$ , and  $u_t = 1$  if not.

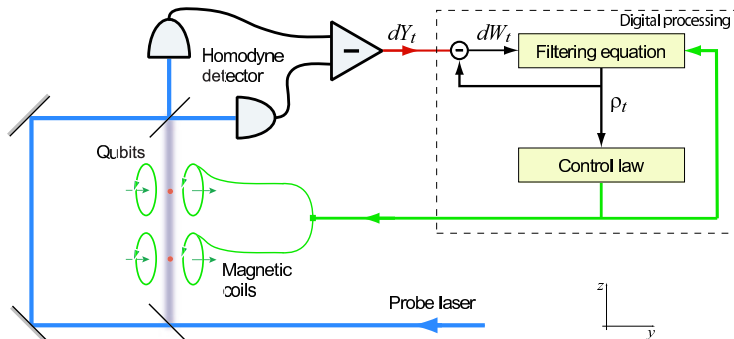
M. Mirrahimi and R. van Handel, SIAM J. Cont. Optimization, 2007.



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# Entangled states: experience scheme



## Example: two qubits

The Hilbert state space for two qubits is given by  $\mathbb{C}^2 \otimes \mathbb{C}^2$ . One can think of an entangled state of the form

$$\frac{1}{\sqrt{2}} (|\uparrow\rangle \otimes |\downarrow\rangle + |\downarrow\rangle \otimes |\uparrow\rangle)$$

# Quantum filter equation

$$d\rho_t = -iu_t^1 [F_y^{(1)}, \rho_t] dt - iu_t^2 [F_y^{(2)}, \rho_t] dt + \frac{1}{2} (2F_z \rho_t F_z - F_z^2 \rho_t - \rho_t F_z^2) dt \\ + (F_z \rho_t + \rho_t F_z - 2\text{Tr}(F_z \rho_t) \rho_t) dW_t, \\ dW_t = dy_t - 2\text{Tr}(F_z \rho_t) dt$$

## 2-qubit system:

- The Hilbert space  $\mathbb{C}^2 \otimes \mathbb{C}^2$  of dimension 4,
- Angular momentum operators

$$F_y^{(1)} = \sigma_y \otimes \text{Id},$$

$$F_y^{(2)} = \text{Id} \otimes \sigma_y,$$

$$F_z = \sigma_z \otimes \text{Id} + \text{Id} \otimes \sigma_z.$$

# Stabilization problem

$$d\rho_t = \frac{1}{2}(2F_z\rho_tF_z - F_z^2\rho_t - \rho_tF_z^2)dt + (F_z\rho_t + \rho_tF_z - 2\text{Tr}(F_z\rho_t)\rho_t)dW_t,$$

## Equilibrium states

$$\Pi_{|\uparrow\uparrow\rangle}, \quad \Pi_{|\downarrow\downarrow\rangle}, \quad \Pi_{\alpha|\uparrow\downarrow\rangle + \beta|\downarrow\uparrow\rangle} \quad (|\alpha|^2 + |\beta|^2 = 1).$$

## Main obstacle

Degeneracy of the measurement operator at the target state: taking the Lyapunov strategy as in the previous case,  $\rho_t$  can converge towards a state in the set  $\{\Pi_{|\uparrow\uparrow\rangle}, \Pi_{|\downarrow\downarrow\rangle}, \Pi_{\alpha|\uparrow\downarrow\rangle + \beta|\downarrow\uparrow\rangle}\}$ .

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# Symmetric and anti-symmetric entangled states

Symmetric entangled state:

$$\rho_S = \Pi \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$

## Idea

We consider a dynamic stabilization by considering:  $u_t^1 = \mathbf{1} + v_t^1$  and  $u_t^2 = -\mathbf{1} + v_t^2$ . Therefore, the uncontrolled system writes:

$$d\rho_t = -i[F_y^{(1)} - F_y^{(2)}, \rho_t]dt + \frac{1}{2}(2F_z\rho_t F_z - F_z^2\rho_t - \rho_t F_z^2)dt + (F_z\rho_t + \rho_t F_z - 2\text{Tr}(F_z\rho_t)\rho_t)dW_t,$$

The only equilibrium state:  $\Pi \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$ .

# Symmetric and anti-symmetric entangled states

## Symmetric case

The feedback law

- 1  $u_t^1 = 1 - \text{Tr} \left( i[F_y^{(1)}, \rho_t] \rho_S \right)$ ,  $u_t^2 = -1 - \text{Tr} \left( i[F_y^{(2)}, \rho_t] \rho_S \right)$  if  $\text{Tr}(\rho \rho_S) \geq \gamma$ ;
- 2  $u_t^1 = 1$ ,  $u_t^2 = 0$  if  $\text{Tr}(\rho \rho_S) \leq \gamma$ ;

stabilize the symmetric entangled state  $\rho_S = \Pi_{\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)}$ .

## Anti-symmetric case

The feedback law

- 1  $u_t^1 = 1 - \text{Tr} \left( i[F_y^{(1)}, \rho_t] \rho_{AS} \right)$ ,  $u_t^2 = 1 - \text{Tr} \left( i[F_y^{(2)}, \rho_t] \rho_{AS} \right)$  if  $\text{Tr}(\rho \rho_{AS}) \geq \gamma$ ;
- 2  $u_t^1 = 1$ ,  $u_t^2 = 0$  if  $\text{Tr}(\rho \rho_{AS}) \leq \gamma$ ;

stabilize the anti-symmetric entangled state  $\rho_{AS} = \Pi_{\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)}$ .

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# General entangled states

$$d\rho_t = -iu_t^1 [F_y^{(1)}, \rho_t] dt - iu_t^2 [F_y^{(2)}, \rho_t] dt + \frac{1}{2} (2F_z \rho_t F_z - F_z^2 \rho_t - \rho_t F_z^2) dt \\ + (F_z \rho_t + \rho_t F_z - 2\text{Tr}(F_z \rho_t) \rho_t) dW_t, \\ dW_t = dy_t - 2\text{Tr}(F_z \rho_t) dt$$

Target state:  $\Pi_{\alpha|\uparrow\downarrow\rangle + \beta|\downarrow\uparrow\rangle}$ .

## Obstacle

We can not remove the degeneracy exactly at our target.

Here we propose a controllability-type result.

# General entangled states

## A simplification

If the system starts at  $\rho_0$  of the form  $\Pi_{\psi_0} = \psi_0\psi_0^*$  with  $\psi_0 \in \mathbb{C}^2$ , the above dynamics become equivalent to:

$$d\psi_t = -i(u_t^1 F_y^{(1)} + u_t^2 F_y^{(2)})\psi_t dt - \frac{1}{2}F_z^2\psi_t dt + F_z\psi_t dW_t,$$
$$\rho_t = \frac{\psi_t\psi_t^*}{\|\psi_t\|^2}.$$

## Idea

By considering exciting control fields, we remove all equilibria except a small neighborhood of the target state and we prove that, the trajectories, almost surely, hit this small neighborhood in finite time.

# Support theorem

## Strook-Varadhan support theorem

Consider the Ito SDE,

$$dx_t = f(x_t)dt + \sigma(x_t)dW_t,$$

and the associated deterministic controlled equation

$$\frac{d}{dt}x_t^u = f(x_t^u) - \frac{1}{2}\nabla\sigma(x_t^u)x_t^u + u(t)\sigma(x_t^u).$$

Consider  $\mathcal{U}$  the set of all piecewise constant functions from  $\mathbb{R}_+$  to  $\mathbb{R}$ , and define

$$\mathcal{S}_x = \overline{\{x^u : u \in \mathcal{U}\}}$$

the set of all controlled trajectories starting at  $x$ . The set  $\mathcal{S}_x$  is the smallest set of the continuous trajectories starting at  $x$  such that

$$\mathbb{P}(\{\omega \in \Omega \mid x_\cdot(\omega) \in \mathcal{S}_x\}) = 1.$$

## Support theorem: application

$$d\psi_t = -i(u_t^1 F_y^{(1)} + u_t^2 F_y^{(2)})\psi_t dt - \frac{1}{2}F_z^2\psi_t dt + F_z\psi_t dW_t,$$

$$\rho_t = \frac{\psi_t\psi_t^*}{\|\psi_t\|^2}.$$

We consider the controlled deterministic equation:

$$d\psi_t = -i(u_t^1 F_y^{(1)} + u_t^2 F_y^{(2)})\psi_t dt - F_z^2\psi_t dt + U_t F_z\psi_t,$$

$$\rho_t = \frac{\psi_t\psi_t^*}{\|\psi_t\|^2}.$$

We show that by taking  $(|\beta_1| \neq |\beta_2|, \left| \frac{\beta_1\beta_2(\alpha_1^2 - \alpha_2^2)}{\alpha_1\alpha_2(\beta_1^2 - \beta_2^2)} \right| \leq \frac{1}{2}, \alpha_1\alpha_2\beta_1\beta_2 < 0)$

$$u_t^1 = \epsilon\beta_1, \quad u_t^2 = \epsilon\beta_2, \quad U_t = 2\frac{\beta_1\beta_2(\alpha_1^2 - \alpha_2^2)}{\alpha_1\alpha_2(\beta_1^2 - \beta_2^2)},$$

the controlled system converges towards a state in an  $\epsilon$ -neighborhood of the target state  $\Pi_{\alpha_1|\uparrow\downarrow\rangle + \alpha_2|\downarrow\uparrow\rangle}$ .

# Stabilization

## Theorem

The following feedback laws ensure the stabilization in an  $\epsilon$ -neighborhood of the target state  $\Pi_{\alpha|\uparrow\downarrow\rangle+\beta|\downarrow\uparrow\rangle}$ :

$$u_t^{(k)} = \begin{cases} -c_k \text{Tr}[i[F_y^{(k)}, \rho_t] \rho_f] & \text{for } V(\rho_t) \leq \epsilon \\ \epsilon \beta_k & \text{for } \epsilon < V(\rho_t) \leq 1 - \delta \\ \gamma_k & \text{for } V(\rho_t) > 1 - \delta \end{cases} \quad k = 1, 2.$$

Where  $\epsilon, c_1, c_2, \beta_1, \beta_2, \gamma_1, \gamma_2 \in \mathbb{R}$ ,  $c_k > 0$ ,  $\epsilon > 0$  sufficiently small and

$$|\beta_1| \neq |\beta_2|, \quad \left| \frac{\beta_1 \beta_2 (\alpha_1^2 - \alpha_2^2)}{\alpha_1 \alpha_2 (\beta_1^2 - \beta_2^2)} \right| \leq 1/2, \quad \alpha_1 \alpha_2 \beta_1 \beta_2 < 0, \quad (\gamma_1, \gamma_2) \neq (\epsilon \beta_1, \epsilon \beta_2).$$

# 10 Random simulations of the 2-qubit system

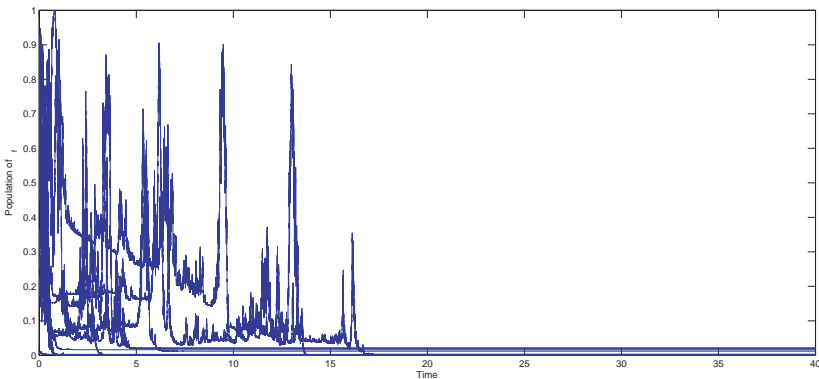


Figure: Approximate stabilization of the .03-neighborhood of  $\frac{1}{\sqrt{5}} |\uparrow\downarrow\rangle + \frac{2}{\sqrt{5}} |\downarrow\uparrow\rangle$