

A curious instability phenomenon for a rounded corner in presence of a negative material

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Coll. with A.-S. Bonnet-Ben Dhia², P. Ciarlet², X. Claeys³ and S.A. Nazarov⁴

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Introduction: general framework

- ▶ Scattering by a **metal** in electromagnetism in **time-harmonic** regime at **optical frequency**.
- ▶ For **metals** at optical frequency, $\Re \varepsilon(\omega) < 0$ and $\Im m \varepsilon(\omega) \ll |\Re \varepsilon(\omega)|$.
⇒ We neglect losses and study the ideal case $\varepsilon(\omega) \in (-\infty; 0)$.

Positive material

$$\varepsilon > 0$$

and $\mu > 0$

Negative metal

$$\varepsilon < 0$$

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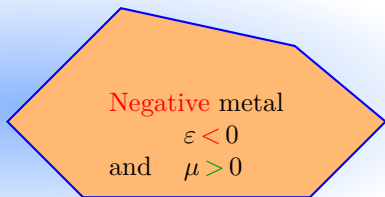
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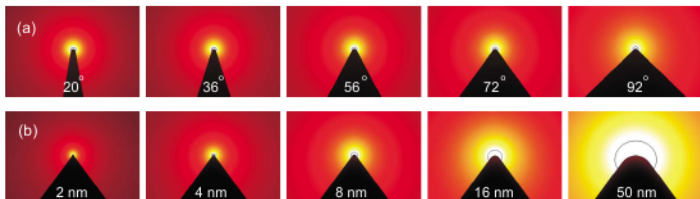
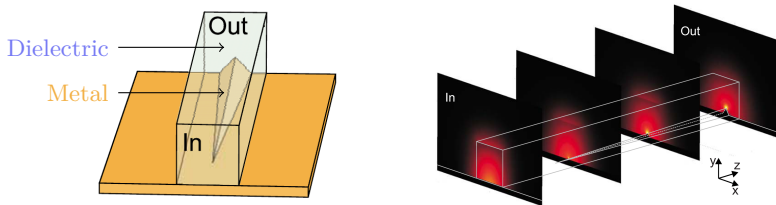
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- ▶ Waves called **Surface Plasmon Polaritons** can propagate **at the interface** between a dielectric and a negative metal.

Introduction: applications

- ▶ **Surface Plasmons Polaritons** can propagate information. Physicists hope to exploit them to reduce the size of **computer chips**.



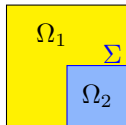
Figures from O'Connor *et al.*, *Appl. Phys. Lett.* 95, 171112 (2009)

- ▶ In this context, physicists use **singular geometries** to **focus energy**. It allows to stock information.

Introduction: in this talk

- ▶ We study a scalar model problem set in a **bounded** domain $\Omega \subset \mathbb{R}^2$:

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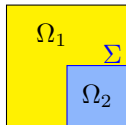


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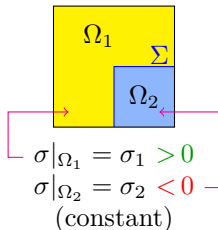


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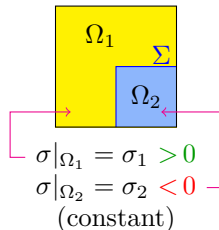


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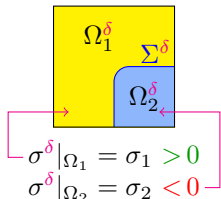
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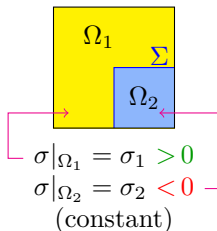
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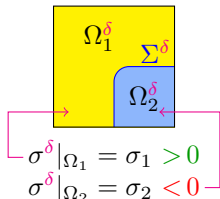
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What is the **behaviour** of the **sequence** $(u^\delta)_\delta$ when δ tends to zero?

Outline of the talk

1 Numerical experiments

To get an **intuition**, we **discretize** (\mathcal{P}^δ) and observe what happens when $\delta \rightarrow 0$.

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We present the properties of the **limit problem** in the geometry with the **real corner** ($\delta = 0$). Since σ changes sign, **original phenomena** appear.

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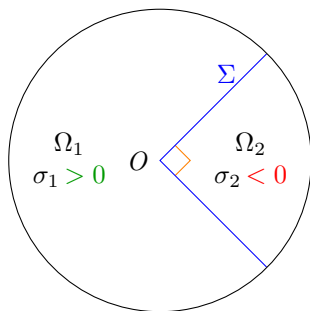
3 Asymptotic analysis

We prove a curious **instability** phenomenon: for certain configurations, (\mathcal{P}^δ) **critically depends** on δ .

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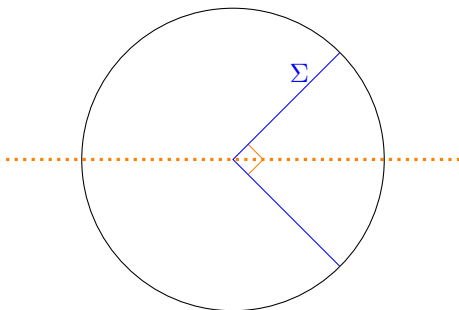
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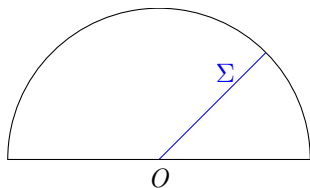
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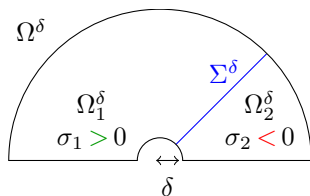
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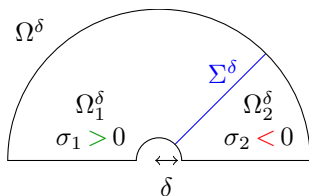
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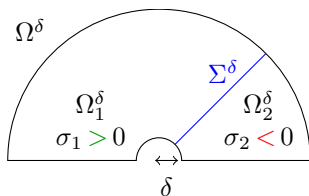
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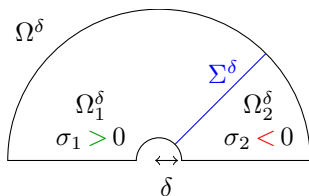
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- ▶ We approximate by a **usual P1 Finite Element Method** this u^δ , *assuming it is well-defined*. The solution of the discretized problem is called u_h^δ .

We display the behaviour of u_h^δ as $\delta \rightarrow 0$.

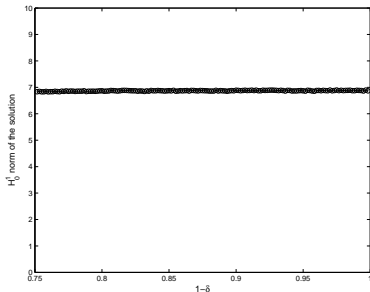
Numerical experiments 1/2

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u_h^δ w.r.t. δ

$\|\nabla u_h^\delta\|_{\Omega^\delta}$ w.r.t. $1 - \delta$

- ▶ For **positive materials**, it is well-known that $(u^\delta)_\delta$ converges to u , the solution in the limit geometry.
- ▶ The **rate of convergence** depends on the **regularity** of u .
- ▶ To avoid to mesh Ω^δ , we can **approximate** u^δ by u_h .

Numerical experiments 2/2

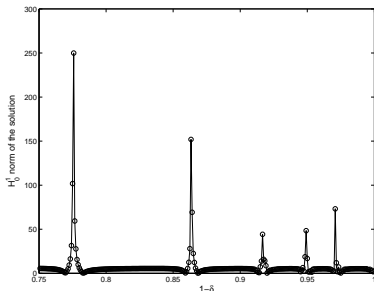
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$$\sigma_1 = 1 \text{ and } \sigma_2 = -0.9999$$

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- For this configuration, u_δ seems to **depend critically** on δ .

In this talk, our goal is to **explain** this behaviour.

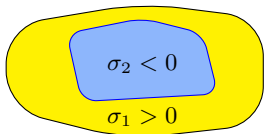
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Problems with a sign changing coefficient

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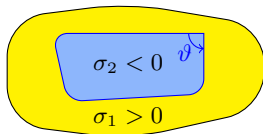
► We have the following properties (see Costabel and Stephan 85, Dauge and Texier 97, Bonnet-Ben Dhia *et al.* 99,10,12,13):

Smooth interface Σ



✓ (\mathcal{P}) well-posed in the Fredholm sense iff $\kappa_\sigma = \sigma_2/\sigma_1 \neq -1$.

Interface Σ with a corner



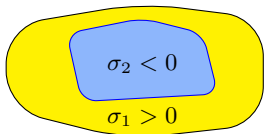
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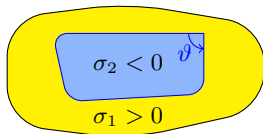
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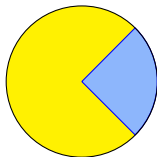


Well-posedness depends on the smoothness of Σ and on σ .

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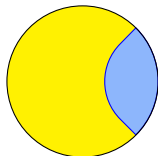
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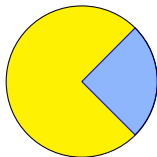
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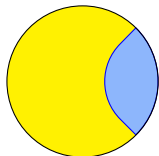
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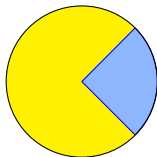
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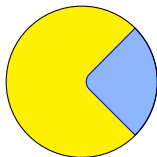
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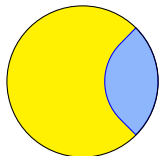
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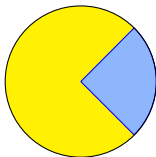
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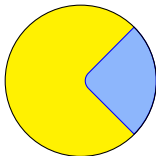
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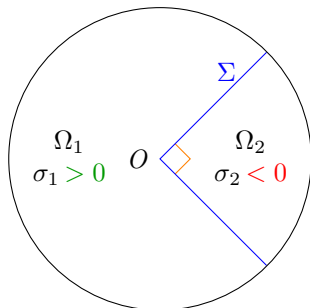


- We need to precise the properties of (\mathcal{P}) when the **interface** has a **corner** in the case $\kappa_\sigma \in I_c \setminus \{-1\}$.

Properties of the limit problem inside the critical interval

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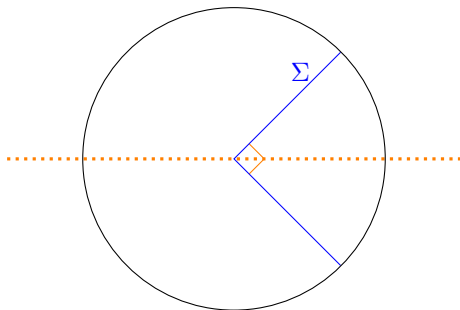
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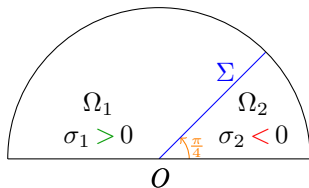
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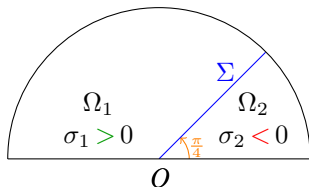
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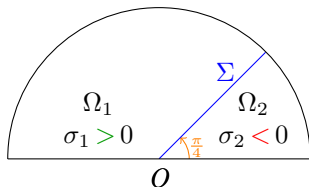
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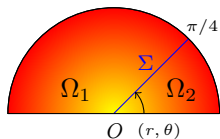
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What happens when $\kappa_\sigma \in (-1; -1/3]$?

Analogy with a waveguide problem

- Bounded sector Ω

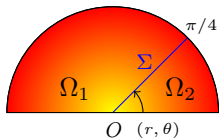


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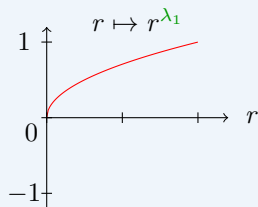
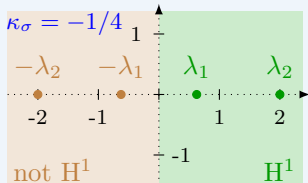
- **Singularities** in the sector

$$s(r, \theta) = r^\lambda \varphi(\theta)$$

Analogy with a waveguide problem

We compute the singularities $s(r, \theta) = r^\lambda \varphi(\theta)$ and we observe two cases:

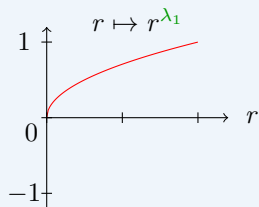
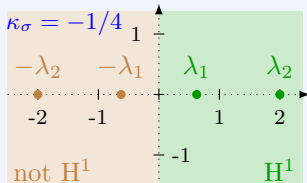
► **Outside the critical interval**



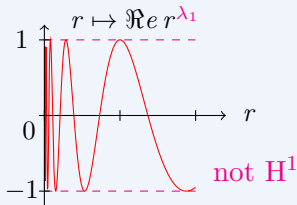
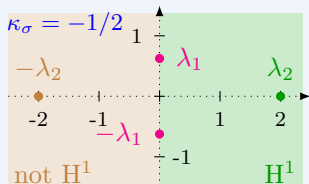
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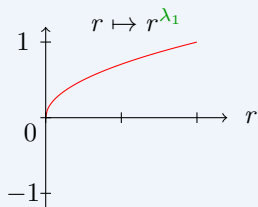
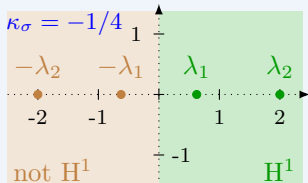
Inside the critical interval



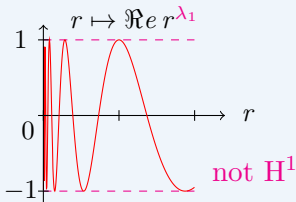
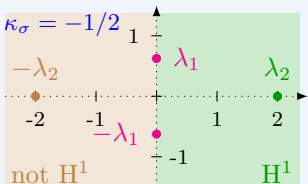
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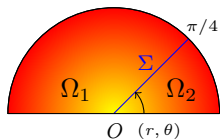
Inside the critical interval



How to deal with the **propagative singularities** inside the critical interval?

Analogy with a waveguide problem

- Bounded sector Ω



- Equation:

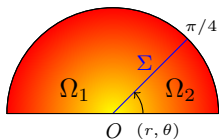
$$\underbrace{-\operatorname{div}(\sigma \nabla u)}_{-r^{-2}(\sigma(r\partial_r)^2 + \partial_\theta \sigma \partial_\theta)u} = f$$

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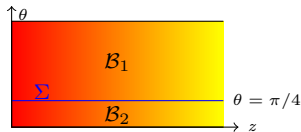
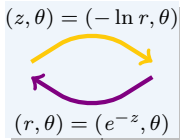
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Analogy with a waveguide problem

- Bounded sector Ω



- Half-strip \mathcal{B}



- Equation:

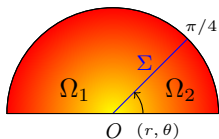
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Analogy with a waveguide problem

- Bounded sector Ω



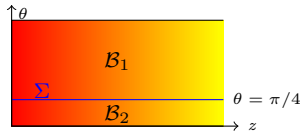
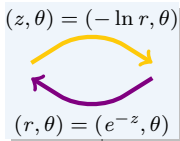
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- Half-strip \mathcal{B}

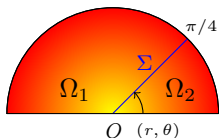


- Equation:

$$\underbrace{-\operatorname{div}(\sigma \nabla u)}_{-(\sigma \partial_z^2 + \partial_\theta \sigma \partial_\theta)u} = e^{-2z} f$$

Analogy with a waveguide problem

- Bounded sector Ω



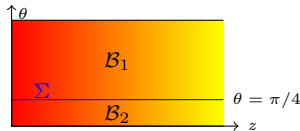
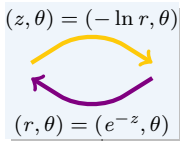
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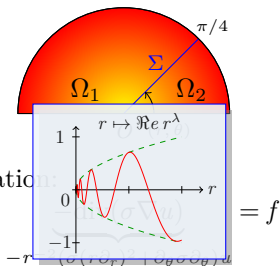
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- Modes in the strip

$$m(z, \theta) = e^{-\lambda z} \varphi(\theta)$$

Analogy with a waveguide problem

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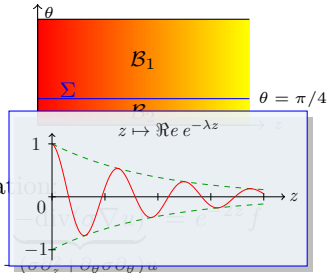
$$s(r, \theta) = r^\lambda \varphi(\theta)$$

$$s \in H^1(\Omega)$$

- Half-strip \mathcal{B}

$$(z, \theta) = (-\ln r, \theta)$$

$$(r, \theta) = (e^{-z}, \theta)$$



- Equation

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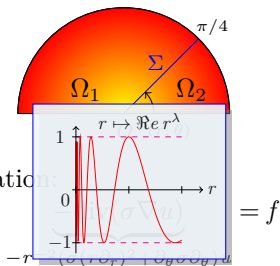
$$m(z, \theta) = e^{-\lambda z} \varphi(\theta)$$

m is evanescent

$$\Re \lambda > 0$$

Analogy with a waveguide problem

- Bounded sector Ω



- Equation:

- Singularities in the sector

$$s(r, \theta) = r^\lambda \varphi(\theta)$$

$$= \cancel{r^a} (\cos b \ln r + i \sin b \ln r) \varphi(\theta)$$

$(\Re \lambda = a, \Im \lambda = b)$

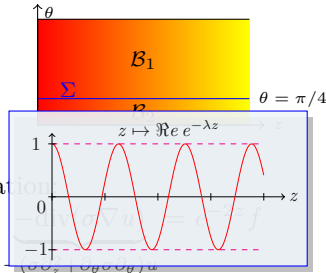
$$s \in H^1(\Omega)$$

$$s \notin H^1(\Omega)$$

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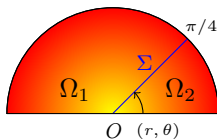
$$m(z, \theta) = e^{-\lambda z} \varphi(\theta)$$

$$= \cancel{e^{-az}} (\cos bz - i \sin bz) \varphi(\theta)$$

m is evanescent
 m is propagative

Analogy with a waveguide problem

- Bounded sector Ω



- Equation:

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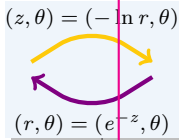
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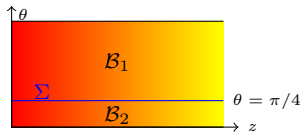
$(\Re \lambda = a, \Im \lambda = b)$

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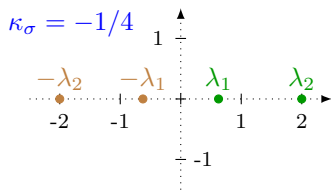
$$\Re \lambda > 0$$

$$\Re \lambda = 0$$

m is **evanescent**
 m is **propagative**

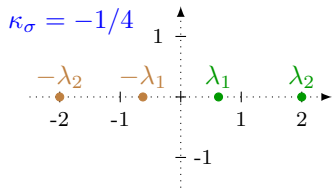
- This encourages us to use **modal decomposition** in the half-strip.

Modal analysis in the waveguide

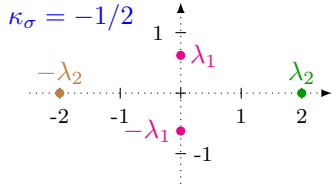


► **Outside the critical interval**. All the modes are exponentially growing or decaying.
→ We look for an exponentially decaying solution. H^1 framework

Modal analysis in the waveguide

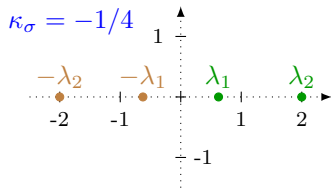


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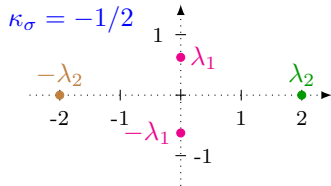


► **Inside the critical interval**. There are exactly two propagative modes.

Modal analysis in the waveguide



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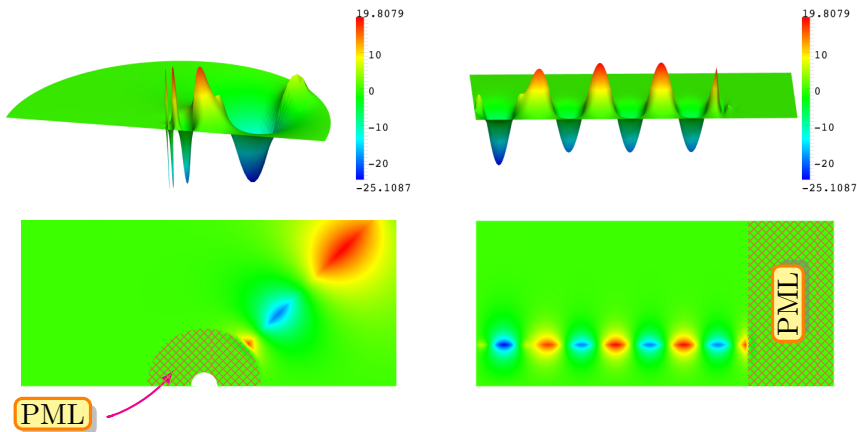
► **Inside the critical interval**. There are exactly two propagative modes.
→ The decomposition on the outgoing modes leads to look for a solution of the form

$$u = \underbrace{c_1 \varphi_1 e^{\lambda_1 z}}_{\text{propagative part}} + \underbrace{u_e}_{\text{evanescent part}}$$

non H^1 framework

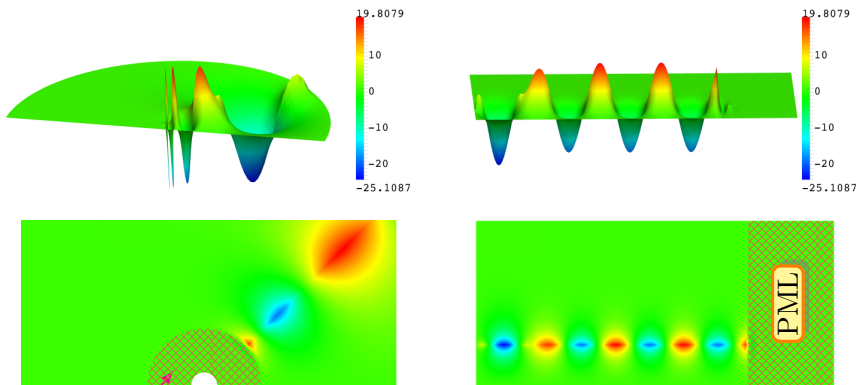
How to approximate the solution?

- ▶ We use a **PML** (*Perfectly Matched Layer*) to bound the domain \mathcal{B} + **finite elements** in the truncated strip ($\kappa_\sigma = -0.999 \in (-1; -1/3)$).



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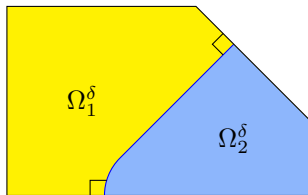
PML



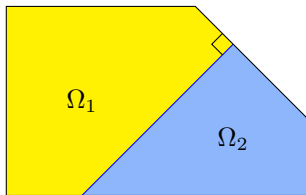
Without the PML, the solution in the **truncated strip** of length L **does not converge** when $L \rightarrow \infty$. This is what we observed in our **numerical experiment** for the **rounded corner**.

- 1 Numerical experiments
- 2 Properties of the limit problem
- 3 Asymptotic analysis**

Asymptotic analysis



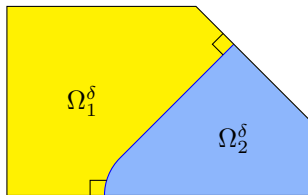
$$(\mathcal{P}^\delta) \quad \left| \begin{array}{l} \text{Find } u^\delta \in H_0^1(\Omega) \text{ s.t.:} \\ -\operatorname{div}(\sigma^\delta \nabla u^\delta) = f \text{ in } \Omega. \end{array} \right.$$



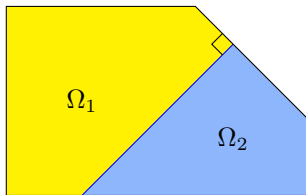
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Asymptotic analysis



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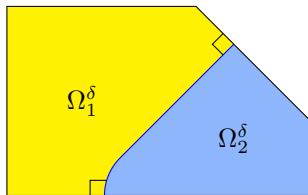


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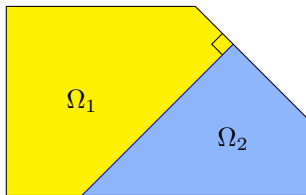
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If (\mathcal{P}) **well-posed** (in $H_0^1(\Omega)$), then u^δ is uniquely defined for δ small enough and $(u^\delta)_\delta$ **converges to u** (as for positive materials).

Asymptotic analysis



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If the limit problem is well-posed only in the **exotic framework**, then (\mathcal{P}^δ) **critically depends** on the value of the **rounding parameter δ** .

Asymptotic analysis

IDEA OF THE APPROACH:

① We prove an *a priori estimate* for u^δ for all δ in some set \mathcal{S} which excludes a discrete set accumulating in zero (♠ hard part of the proof, S.A. Nazarov's technique).



$$\ln \mathcal{S} = \{\ln \delta, \delta \in \mathcal{S}\}$$

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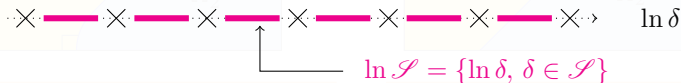
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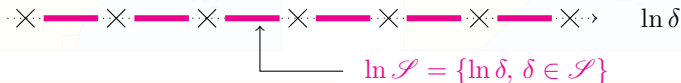
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$$\|u^\delta - \hat{u}^\delta\|_{H_0^1(\Omega^\delta)} \leq c \delta^\beta \|f\|_{\Omega^\delta}, \quad \forall \delta \in \mathcal{S}.$$

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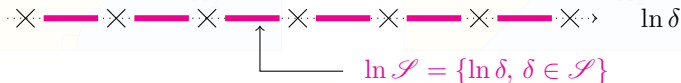
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③ The behaviour of $(\hat{u}^\delta)_\delta$ can be explicitly examined as $\delta \rightarrow 0$. **The sequence $(\hat{u}^\delta)_\delta$ does not converge, even for the L^2 -norm!**

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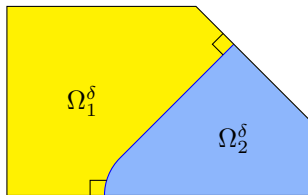
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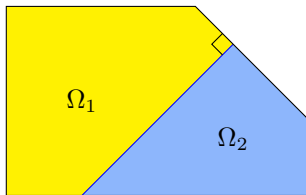
④ Conclusion.

The sequence $(u^\delta)_\delta$ **does not converge**, even for the L^2 -norm!

Asymptotic analysis



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- The behaviour of $(u^\delta)_\delta$ depends on the properties of the **limit problem**.

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Conclusion

Let us remind the initial question:



What is the **behaviour** of $(u^\delta)_\delta$ when δ tends to zero?

Conclusion

Let us remind the initial question:



What is the **behaviour** of $(u^\delta)_\delta$ when δ tends to zero?



This depends on the features of the **limit problem**.

(...)

(...)

$$\kappa_\sigma = -1.0001 \notin I_c$$

$$\kappa_\sigma = -0.9999 \in I_c$$



When $\kappa_\sigma \in I_c$, $(u^\delta)_\delta$ **does not converge**, even for the L^2 -norm!

In this case, it is impossible to **simulate** the fields since it is impossible to **measure** exactly δ . \Rightarrow What happens **physically**?

Thank you for your attention!!!



A.-S. Bonnet-Ben Dhia, L. Chesnel, P. Ciarlet Jr., *T-coercivity for scalar interface problems between dielectrics and metamaterials*, *M2AN*, 46, 1363–1387, 2012.



A.-S. Bonnet-Ben Dhia, L. Chesnel, X. Claeys, *Radiation condition for a non-smooth interface between a dielectric and a metamaterial*, *M3AS*, 23, 2013.



L. Chesnel, X. Claeys, S.A. Nazarov, *A curious instability phenomenon for a rounded corner in presence of a negative material*, [preprint arXiv:1304.4788](https://arxiv.org/abs/1304.4788), 2013.

Naive approximation

- ▶ Let us try a **usual Finite Element Method** (P1 Lagrange Finite Element). We solve the problem

$$\left| \begin{array}{l} \text{Find } u_h \in V_h \text{ s.t.:} \\ \int_{\Omega} \sigma \nabla u_h \cdot \nabla v_h = \int_{\Omega} f v_h, \quad \forall v \in V_h, \end{array} \right.$$

where V_h approximates $H_0^1(\Omega)$ as $h \rightarrow 0$ (h is the **mesh size**).

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THE SEQUENCE (u_h) DOES NOT CONVERGE AS $h \rightarrow 0$!!!

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(...)

(...)

Contrast $\kappa_{\sigma} = -0.999 \in (-1; -1/3)$.

Naive approximation

- ▶ Outside the critical interval, the sequence (u_h) converges.

(...)

(...)

Contrast $\kappa_\sigma = -1.001 \notin (-1; -1/3)$.